

Karhunen-Loève Expansions of Lévy Processes

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1 Introduction

2 Lévy Processes and Infinitely Divisible Random Vectors

3 Main Results

- KLE Components
- Simulation

4 Examples

We want to expand a continuous time stochastic process X in a stochastic Fourier series on the interval $[0, T]$:

$$X_t = \sum_{k \geq 1} Y_k \phi_k(t),$$

where $\{\phi_k\}_{k \geq 1}$ is an orthonormal basis of $L^2([0, T], \mathbb{R})$, and our stochastic Fourier coefficients are given by

$$Y_k := \int_0^T X_t \phi_k(t) dt.$$

Things to think about:

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- ▶ How should we choose $\{\phi_k\}_{k \geq 1}$?
- ▶ Distribution of, and dependence among $\{Y_k\}_{k \geq 1}$?

Assumptions: (a) $\mathbb{E}[X_t] = 0$, (b) $\mathbb{E}[X_t^2] < \infty$, and (c) $\text{Cov}(X_s, X_t)$ is continuous

Basis: The eigenfunctions $\{e_k\}_{k \geq 1}$ corresponding to the non-zero eigenvalues $\{\lambda_k\}_{k \geq 1}$ of the operator $K : L^2([0, T]) \rightarrow L^2([0, T])$,

$$(Kf)(s) := \int_0^T \text{Cov}(X_s, X_t) f(t) dt$$

constitute a basis for $L^2([0, T])$. We define

$$Z_k := \int_0^T X_t e_k(t) dt$$

and determine the order of $\{e_k\}_{k \geq 1}$, $\{Z_k\}_{k \geq 1}$, and $\{\lambda_k\}_{k \geq 1}$ according to $\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq \dots$

Theorem (The Karhunen-Loève Theorem (KLT))

(i)

$$\mathbb{E} \left[\left(X_t - \sum_{k=1}^d Z_k e_k(t) \right)^2 \right] \rightarrow 0, \quad \text{as } d \rightarrow \infty$$

uniformly for $t \in [0, T]$. Additionally, the $\{Z_k\}_{k \geq 1}$ are uncorrelated and satisfy $\mathbb{E}[Z_k] = 0$ and $\mathbb{E}[Z_k^2] = \lambda_k$.

(ii) For any other basis $\{\phi_k\}_{k \geq 1}$ with corresponding Fourier coefficients $\{Y_k\}_{k \geq 1}$, and any $d \in \mathbb{N}$, we have

$$\int_0^T \mathbb{E} \left[(\varepsilon_d(t))^2 \right] dt \leq \int_0^T \mathbb{E} \left[(\tilde{\varepsilon}_d(t))^2 \right] dt,$$

where ε_d and $\tilde{\varepsilon}_d$ are the remainders $\varepsilon_d(t) := \sum_{d+1}^{\infty} Z_k e_k(t)$ and $\tilde{\varepsilon}_d(t) := \sum_{d+1}^{\infty} Y_k \phi_k(t)$.

Proof in:



GHANEM, R. G.. AND SPANOS, P.D.. (1991).
Stochastic finite elements: A spectral approach.
Springer-Verlag, New York-Berlin-Heidelberg.

They credit:



KAC, M. AND SIEGERT, A. (1947).

An explicit representation of a stationary Gaussian process.
Ann. Math. Stat. **18**, 438–442.



KARHUNEN, K. (1947).

Über lineare Methoden in der Wahrscheinlichkeitsrechnung.
Amer. Acad. Sc. Fennicade, Ser. A, I **37**, 3–79.



LOÉVE, M. (1948).

Fonctions aleatoires du second ordre.

In *Processus stochastic et mouvement Brownien*. ed. P. Lévy.
Gauthier Villars, Paris.

In order to use the KLT for a chosen process X in any meaningful way we see that we need to determine:

- ▶ The eigenvalues $\{\lambda_k\}_{k \geq 1}$ and eigenfunctions $\{e_k\}_{k \geq 1}$ of the operator K . I.e. solve:

$$\int_0^T \text{Cov}(X_s, X_t) e_k(s) ds = \lambda_k e_k(t),$$

a homogeneous Fredholm integral equation of the second kind

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- ▶ The distribution of the Fourier coefficients $\{Z_k\}_{k \geq 1}$
- ▶ If we want to simulate an approximate path of X via a Karhunen Loève expansion (KLE) we also need to know how to simulate the first d Fourier coefficients, i.e. simulate the random vector

$$Z^{(d)} = (Z_1, \dots, Z_d).$$

Note: The components of this vector need not be independent!

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BARCZY, M. AND LOVAS, R.L.. (2016).

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PHOON, K.K., HUANG, H.W.. AND QUEK, S.T.. (2005).

Simulation of strongly non-Gaussian processes using Karhunen–Loeve expansion.

Probabilistic Engineering Mechanics **20**, 188–198.

Goal: For a one-dimensional Lévy process X satisfying $\mathbb{E}[X_t] = 0$ and $\mathbb{E}[X_t^2] < \infty$ determine

- ▶ $\{\lambda_k\}_{k \geq 1}$, $\{e_k\}_{k \geq 1}$

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- ▶ The distribution and dependence structure of $Z^{(d)} = (Z_1, \dots, Z_d)$
- ▶ How to simulate $Z^{(d)}$

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- ▶ The distribution of a d -dimensional Lévy process X or infinitely divisible vector ξ is completely determined by the *characteristic exponent*:

$$\Psi_X(\mathbf{z}) := -\frac{1}{t} \log \mathbb{E}[e^{i\langle \mathbf{z}, X_t \rangle}], \quad \text{or} \quad \Psi_\xi(\mathbf{z}) := -\log \mathbb{E}[e^{i\langle \mathbf{z}, \xi \rangle}].$$

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- ▶ The characteristic exponent always has the form

$$\Psi(\mathbf{z}) = \frac{1}{2} \mathbf{z}^\top Q \mathbf{z} - i \langle \mathbf{a}, \mathbf{z} \rangle - \int_{\mathbb{R}^n \setminus \{\mathbf{0}\}} e^{i\langle \mathbf{z}, \mathbf{x} \rangle} - 1 - i \langle \mathbf{z}, \mathbf{x} \rangle h(\mathbf{x}) \nu(d\mathbf{x}).$$

So Ψ and therefore the distribution are determined by the generating triple $(\mathbf{a}, Q, \nu)_h$ where

- ▶ $\mathbf{a} \in \mathbb{R}^d$ ($a \in \mathbb{R}$),
- ▶ $Q \in \mathbb{R}^{d \times d}$ is positive semi definite ($\sigma^2 \geq 0$) (*Gaussian component*),
and
- ▶ ν is a measure (*Lévy measure*)
- ▶ h a *cut-off* function needed in general to make the integral converge

If

$$\int_{|\mathbf{x}| < 1} |\mathbf{x}| \nu(d\mathbf{x}) < \infty$$

then we can set $h \equiv 0$. (X has bounded variation when $Q = 0$ or $\sigma^2 = 0$.)

A scaled Brownian motion with drift with $\mathbb{E}[X_t] = \mu t$ and $\text{Var}(X_t) = \sigma^2 t$ has characteristic exponent

$$\Psi(z) = \frac{\sigma^2}{2} z^2 - i\mu z.$$

Since $\nu \equiv 0$ there are no jumps and we have generating triple $(\mu, \sigma^2, 0)$

For $c, \rho, \hat{\rho} > 0$ set

$$\nu(\mathrm{d}x) = \mathbb{I}(x < 0) c \frac{e^{\hat{\rho}x}}{|x|} \mathrm{d}x + \mathbb{I}(x > 0) c \frac{e^{-\rho x}}{x} \mathrm{d}x.$$

Then

$$\Psi(z) = - \int_{\mathbb{R} \setminus \{0\}} \left(e^{izx} - 1 \right) \nu(\mathrm{d}x) = c \left(\log \left(1 - \frac{iz}{\rho} \right) + \log \left(1 + \frac{iz}{\hat{\rho}} \right) \right)$$

is the characteristic exponent of a variance gamma process with generating triple $(0, 0, \nu)$.

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Under the KLT assumption it is straight-forward to show that

$$\text{Var}(X_t) = \mathbb{E}[X_t^2] = \alpha t, \quad \text{and} \quad \text{Cov}(X_s, X_t) = \alpha \min(s, t),$$

where $\alpha = \Psi''(0)$.

For example, for our Brownian motion we have $\alpha = \sigma^2$ and for our variance gamma process $\alpha = c(\rho^{-2} + \hat{\rho}^{-2})$.

But this shows that the eigenvalues/functions of the operator K for a Lévy process can be derived in exactly the same manner as those of a Brownian motion. And so we get our first result essentially for free...

Proposition

The eigenvalues and associated eigenfunctions of the operator K defined with respect to a Lévy process X are given by


$$\lambda_k = \frac{\alpha T^2}{\pi^2 \left(k - \frac{1}{2}\right)^2}, \quad \text{and} \quad e_k(t) = \sqrt{\frac{2}{T}} \sin\left(\frac{\pi}{T} \left(k - \frac{1}{2}\right) t\right),$$

for $k \in \mathbb{N}$ and $t \in [0, T]$.

Idea of the Proof: Differentiate both sides of

$$\alpha \int_0^T \min(s, t) e_k(s) ds = \lambda_k e_k(t),$$

twice with respect to t to reduce to an ODE. Details for Brownian Motion case in

-  [ASH, R.B.. AND GARDNER, M.F.. \(1975\).](#)
Topics in stochastic process.
Academic Press, New York–San Francisco–London.

Define the *total variance* of X on $[0, T]$ as

$$v(T) := \int_0^T \mathbb{E}[X_t^2] dt = \frac{\alpha T^2}{2}$$

But we also have $v(T) = \sum_{k \geq 1} \lambda_k$ (in general). Therefore, the total variance explained by a d -term approximation is

$$\frac{\sum_{k=1}^d \lambda_k}{v(T)} = \frac{2}{\pi^2} \sum_{k=1}^d \frac{1}{\left(k - \frac{1}{2}\right)^2}.$$

Computation yields: the first 2, 5 and 21 terms already explain 90%, 95%, and 99% of the total variance of the process. Holds independently of X , α or T .

A key property of Gaussian processes, is that the coefficients $\{Z_k\}_{k \geq 1}$ are again Gaussian. For example,

$$B_t = \sqrt{2} \sum_{k \geq 1} Z_k \frac{\sin\left(\pi\left(k - \frac{1}{2}t\right)\right)}{\pi\left(k - \frac{1}{2}\right)}$$

where B is a standard Brownian Motion on $[0, 1]$ and the $\{Z_k\}_{k \geq 1}$ are i.i.d. random variables with common distribution $\mathcal{N}(0, 1)$.

Lemma

Let X be a Lévy process and $\{f_k\}_{k=1}^d$ be a collection of functions which are in $L^1([0, T])$. Then the vector ξ consisting of elements

$$\xi_k = \int_0^T X_t f_k(s) ds, \quad k \in \{1, 2, \dots, d\},$$

has an ID distribution with characteristic exponent

$$\Psi_\xi(\mathbf{z}) = \int_0^T \Psi_X(\langle \mathbf{z}, \mathbf{u}(t) \rangle) dt, \quad \mathbf{z} \in \mathbb{R}^d,$$

where $\mathbf{u} : [0, T] \rightarrow \mathbb{R}^d$ is the function with k -th component $u_k(t) := \int_t^T f_k(s) ds$, $k \in \{1, 2, \dots, d\}$.

Theorem

If X is a Lévy process with generating triple (a, σ^2, ν) that satisfies the KLT and BV assumptions then $Z^{(d)}$ has generating triple $(\mathbf{a}, \mathcal{Q}, \Pi)$ where \mathbf{a} is the vector with entries

$$a_k := a \frac{(-1)^{k+1} \sqrt{2} T^{\frac{3}{2}}}{\pi^2 \left(k - \frac{1}{2}\right)^2}, \quad k \in \{1, 2, \dots, d\},$$

Theorem (cont.)

Q is a diagonal $d \times d$ matrix with entries

$$q_{k,k} := \frac{\sigma^2}{2} \frac{T^2}{\pi^2 \left(k - \frac{1}{2}\right)^2}, \quad k \in \{1, 2, \dots, d\},$$

and Π is the measure,

$$\Pi(B) := \int_{\mathbb{R} \setminus \{0\} \times [0, T]} \mathbb{I}(f(\mathbf{v}) \in B) (\nu \times \lambda)(d\mathbf{v}), \quad B \in \mathcal{B}_{\mathbb{R}^d \setminus \{0\}},$$

where λ is the Lebesgue measure on $[0, T]$ and $f : \mathbb{R} \times [0, T] \rightarrow \mathbb{R}^d$ is the function

$$(x, t) \mapsto \frac{\sqrt{2T}x}{\pi} \left(\frac{\cos\left(\frac{\pi}{T}\left(1 - \frac{1}{2}\right)t\right)}{\left(1 - \frac{1}{2}\right)}, \dots, \frac{\cos\left(\frac{\pi}{T}\left(d - \frac{1}{2}\right)t\right)}{\left(d - \frac{1}{2}\right)} \right)^{\mathbf{T}}.$$

Idea of Proof:

- ▶ From Lemma: $\Psi_{Z^{(d)}}(z) = \int_0^T \Psi_X(\langle \mathbf{z}, \mathbf{u}(t) \rangle) dt$

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$$u_k(t) = \sqrt{\frac{2}{T}} \int_t^T \sin\left(\frac{\pi}{T} \left(k - \frac{1}{2}\right) s\right) ds = \sqrt{2T} \frac{\cos\left(\frac{\pi}{T} \left(k - \frac{1}{2}\right) t\right)}{\pi\left(k - \frac{1}{2}\right)}$$

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- ▶ Evaluation of integral over $[0, T]$ and the orthogonality of $\{u_k\}_{1 \leq k \leq d}$ give \mathbf{a} and \mathcal{Q} .

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- ▶ From Lemma: $\Psi_{Z^{(d)}}(z) = \int_0^T \Psi_X(\langle \mathbf{z}, \mathbf{u}(t) \rangle) dt$
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$$u_k(t) = \sqrt{\frac{2}{T}} \int_t^T \sin\left(\frac{\pi}{T} \left(k - \frac{1}{2}\right) s\right) ds = \sqrt{2T} \frac{\cos\left(\frac{\pi}{T} \left(k - \frac{1}{2}\right) t\right)}{\pi\left(k - \frac{1}{2}\right)}$$

- ▶ Evaluation of integral over $[0, T]$ and the orthogonality of $\{u_k\}_{1 \leq k \leq d}$ give \mathbf{a} and \mathcal{Q} .
- ▶ Fubini's theorem and a change of variables yields Π .

Corollary

$Z^{(d)}$ has independent entries if, and only if, ν is the zero measure.

Idea of Proof:

- ▶ (\Leftarrow) known, or because Q is diagonal
- ▶ (\Rightarrow) we show that Π is not supported by the coordinate axes

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How can we simulate a random vector with dependent entries with only knowledge of the characteristic function?

In general, this seems to be a difficult problem. Infinite divisibility makes it possible!

Idea: Write $Z^{(d)}$ as a(n) (infinite) sum of simpler random vectors.

- ▶ Random sequences $\{V_i\}_{i \geq 1}$ and $\{\Gamma_i\}_{i \geq 1}$ which are independent of each other and defined on a common probability space.

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- ▶ $\{V_i\}_{i \geq 1}$ are independent and take values in some measurable space D with common dist. F
- ▶ $H : (0, \infty) \times D \rightarrow \mathbb{R}^d$, measurable, and

$$S_n := \sum_{i=1}^n H(\Gamma_i, V_i), \quad n \in \mathbb{N},$$

and

$$\mu(B) := \int_0^\infty \int_D \mathbb{I}(H(r, v) \in B) F(\mathrm{d}v) \mathrm{d}r, \quad B \in \mathcal{B}_{\mathbb{R}^d \setminus \{0\}}.$$

Theorem

If μ is a Lévy measure satisfying the BV assumption, then S_n converges almost surely to an ID random vector with generating triple $(\mathbf{0}, \mathbf{0}, \mu)$ as $n \rightarrow \infty$.

Theorems 3.1, 3.2, and 3.4 in



ROSIŃSKI, J. (1990).

On series representations of infinitely divisible random vectors.

The Annals of Probability **18**, 405–430.

Idea: Show that Π has a disintegrated form like μ .

Simplifying assumptions:

- ▶ X has no Gaussian component
- ▶ X has only positive jumps
- ▶ $X_t \stackrel{d}{=} X_t^+ - X_t^- + B_t \Rightarrow Z_X^{(d)} \stackrel{d}{=} Z_{X^+}^{(d)} - Z_{X^-}^{(d)} + Z_B^{(d)}$

Necessary assumption: ν has a strictly positive density π .

We want

$$g(x) := \int_x^\infty \pi(s) ds$$

to have a well defined, non-increasing inverse g^{-1} .

Theorem

If X is a Lévy process with only positive jumps, satisfying the KLT and BV assumptions, such that X has triple $(a, 0, \pi)$, where π is strictly positive, then

$$Z^{(d)} \stackrel{d}{=} \mathbf{a} + \sum_{i \geq 1} H(\Gamma_i, U_i),$$

where $\{U_i\}_{i \geq 1}$ is an i.i.d. sequence of uniform random variables on $[0, 1]$, and

$$H(r, v) := f(g^{-1}(r/T), Tv),$$

where f is the function defined previously, i.e.

$$f(x, t) = \frac{\sqrt{2T}x}{\pi} \left(\frac{\cos\left(\frac{\pi}{T}\left(1 - \frac{1}{2}\right)t\right)}{\left(1 - \frac{1}{2}\right)}, \dots, \frac{\cos\left(\frac{\pi}{T}\left(d - \frac{1}{2}\right)t\right)}{\left(d - \frac{1}{2}\right)} \right)^{\mathbf{T}}.$$

$$S_t^{(d)} := \sqrt{\frac{2}{T}} \sum_{k=1}^d Z_k \sin \left(\frac{\pi}{T} \left(k - \frac{1}{2} \right) t \right)$$

$$Z_k \stackrel{d}{=} a_k + \sum_{i \geq 1} \frac{\sqrt{2T} g^{-1}(\Gamma_i/T) \cos \left(\pi \left(k - \frac{1}{2} \right) U_i \right)}{\pi \left(k - \frac{1}{2} \right)},$$

Nice/Unique features:

- ▶ $Z^{(d)}$ is independent of t

Potentially difficult:

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- ▶ Each summand of $\sum_{i \geq 1} H(\Gamma_i, U_i)$ requires the evaluation of d cosine functions

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- ▶ Each summand of $\sum_{i \geq 1} H(\Gamma_i, U_i)$ requires the evaluation of d cosine functions
- ▶ Inverting g

$$S_t^{(d)} := \sqrt{\frac{2}{T}} \sum_{k=1}^d Z_k \sin \left(\frac{\pi}{T} \left(k - \frac{1}{2} \right) t \right)$$

$$Z_k \stackrel{d}{=} a_k + \sum_{i \geq 1} \frac{\sqrt{2T} g^{-1}(\Gamma_i/T) \cos \left(\pi \left(k - \frac{1}{2} \right) U_i \right)}{\pi \left(k - \frac{1}{2} \right)},$$

Nice/Unique features:

- ▶ $Z^{(d)}$ is independent of t
- ▶ We can increase d incrementally very easily (without starting fresh simulation)
- ▶ $S^{(d)}$ has smooth paths

Potentially difficult:

- ▶ Each summand of $\sum_{i \geq 1} H(\Gamma_i, U_i)$ requires the evaluation of d cosine functions
- ▶ Inverting g
- ▶ Convergence of sum depends on decay of g^{-1}

- 1 Introduction
- 2 Lévy Processes and Infinitely Divisible Random Vectors
- 3 Main Results
 - KLE Components
 - Simulation
- 4 Examples

Consider a VG process with $c = 1$, $\rho = 2$, $\hat{\rho} = 5$, and $T = 3$. The process X^+ has Lévy measure $\nu(dx) = \frac{e^{-\rho x}}{x} dx$ so that we have

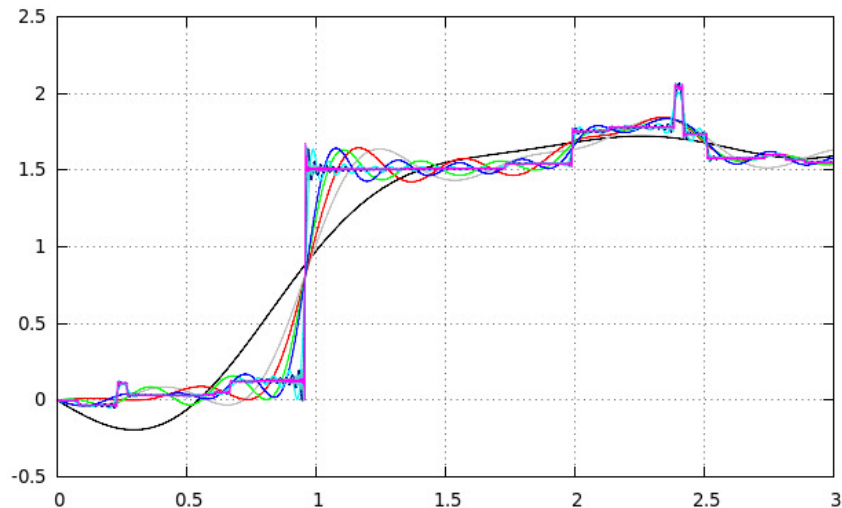
$$g(x) = c \int_x^\infty \frac{e^{-\rho x}}{x} dx = cE_1(\rho x), \quad \text{and} \quad g^{-1}\left(\frac{\Gamma_i}{T}\right) = \frac{1}{\rho} E_1^{-1}\left(\frac{\Gamma_i}{Tc}\right),$$

where E_1 is the exponential integral function.

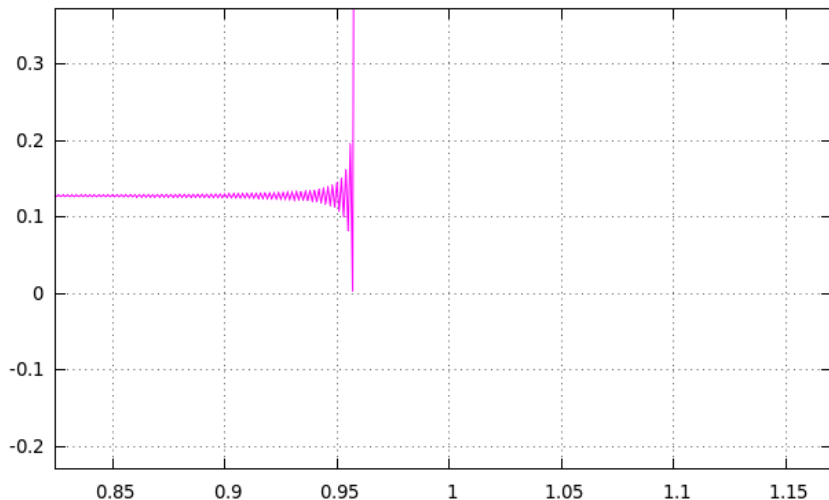
We truncate the random series when $\Gamma_i/(Tc) > 46$ since at this point $g^{-1}(\Gamma_i/T) < \rho^{-1}10^{-19}$. We expect to generate $46Tc = 138$ random pairs (U_i, Γ_i) per path of X^+ .

Example 1

Sample paths of $S^{(d)}$ with $d \in \{5, 10, 15, 20, 25, 100, 250, 500, 3000\}$

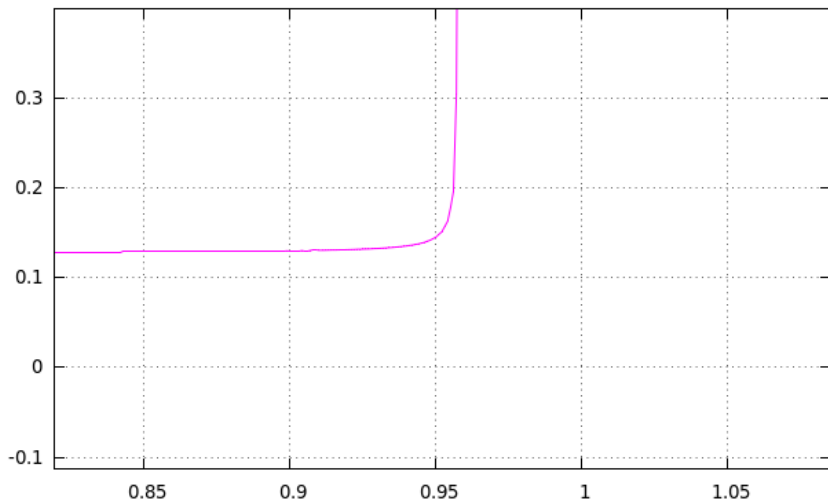


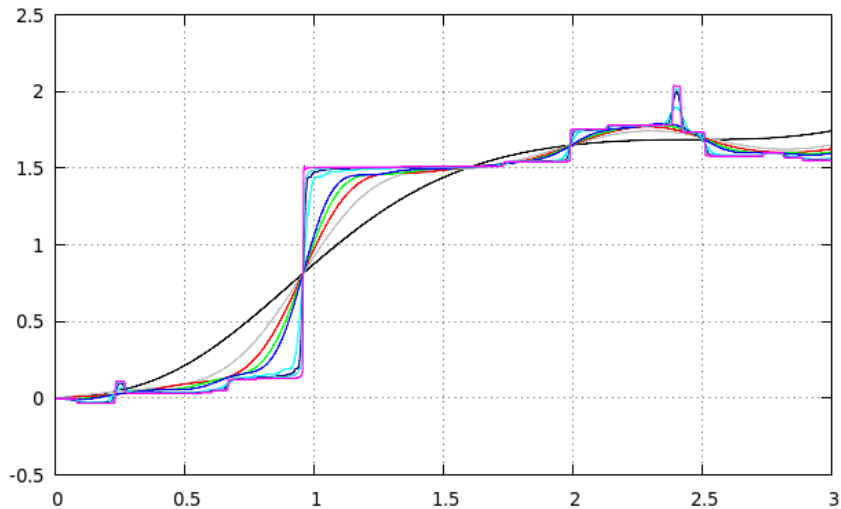
A closer look at $d = 3000$ shows...



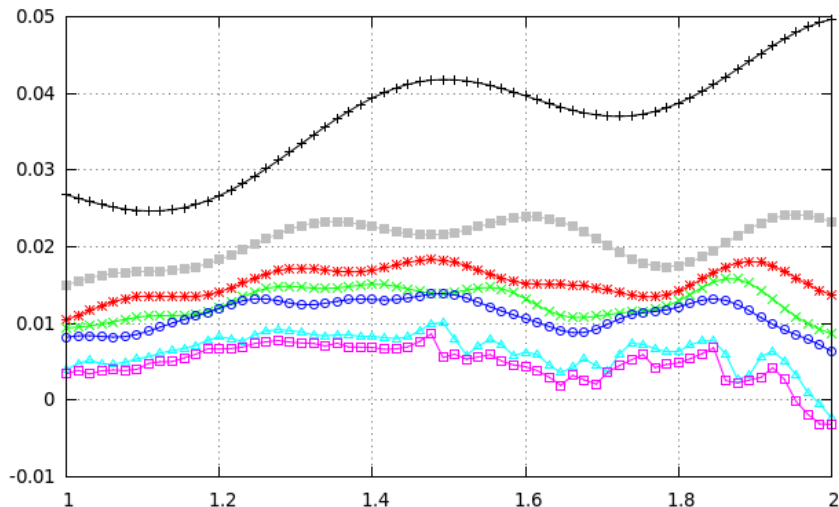
... the Gibbs phenomenon...

... but we can remove this if we want. For example if we form the Cesàro sums $C_t^{(d)} := \frac{1}{d} \sum_{k=1}^d S_t^{(k)}$ the Gibbs phenomenon disappears.





Suppose we want to compute $\mathbb{E}[e^{X_t}] = e^{-\Psi(-i)} = (5/3)^t$ on the interval $[1, 2]$ by simulating 10^6 paths of $S^{(d)}$. The errors:



This method could be useful if we have to evaluate $\mathbb{E}[f(X_t)]$ at many different points in the interval $[0, T]$, or at previously unknown points, say if we wanted to find the minimum in t .

If, for example, $d \leq 100$ is deemed good enough, then 100×10^6 single precision floating point numbers (realizations of the $\{Z_k\}_{1 \leq k \leq d}$) can be stored in memory: $4 \text{ Bytes} \times 10^8 \approx 0.37 \text{ GB}$. Even $d = 3000$ is not impossible: $4 \text{ Bytes} \times 3 \times 10^9 \approx 11.2 \text{ GB}$.

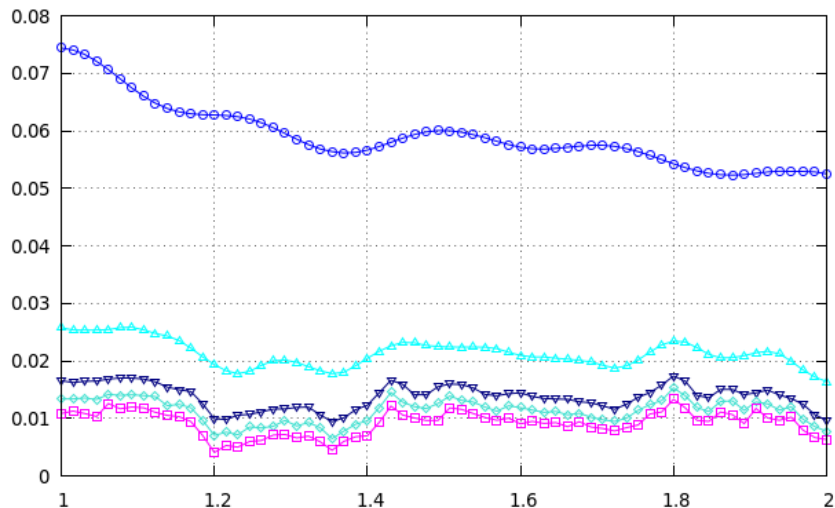
If we want to compute $\mathbb{E}[f(X_t)]$ at points $t_1 \leq t_2 \leq \dots \leq t_d$ in $[0, T]$ then we can also use a fractional FFT algorithm to compute the realizations of $\{S_{t_k}^{(d)}\}_{1 \leq k \leq d}$ with $O(d \log(d))$ operations.

Consider a VG process X with $c = 5$, $\rho = 21.8735$, $\hat{\rho} = 56.4414$, and $T = 3$. We will also add a drift μ whose value is to be determined. Suppose we want to compute

$$e^{-r\tau} \mathbb{E} \left[(A_0 \exp(X_\tau) - K)^+ \right], \quad \text{or}$$
$$e^{-r\tau} \mathbb{E} \left[\left(\frac{A_0}{\tau} \int_0^\tau \exp(X_t) dt - K \right)^+ \right],$$

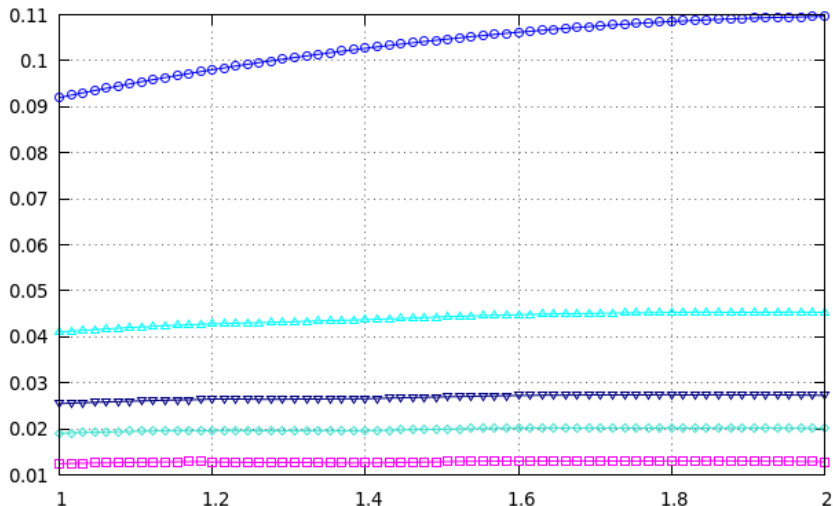
where $r, A_0, K > 0$. Then μ is determined by the risk neutral condition $\Psi(-i) = -r$, and the expressions represent the price of a European and Asian call option respectively.

Replace X by $S^{(d)}$ and do Monte Carlo with 10^6 iterations. Errors for $\tau \in [1, 2]$ compared to a benchmark computed using a Fourier transform technique.



Asian option

Replace X by $C^{(d)}$ and do Monte Carlo with 10^6 iterations. Errors for $\tau \in [1, 2]$ compared to a benchmark computed using a Fourier transform technique.



Let $c^{(d)}$ be a realization of $C^{(d)}$ on $[0, T]$. Then $t \mapsto \exp(c_t^{(d)})$ is smooth function of t . Applying the right change of variables $\cos(\theta) = 2\left(\frac{t}{\tau} - 1\right)$ transforms


$$\int_0^\tau \exp\left(c_t^{(d)}\right) dt$$


into the integral of a periodic function over $[0, \pi]$. We expect exponential convergence of the trapezoidal rule for such a function. Numerical evaluation is then possible via the Clenshaw-Curtis quadrature.


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- ▶  HACKMANN, D. (2016).
Karhunen–Loève expansions of Lévy processes. preprint: arXiv:1603.00677.